Introduction

Understanding the bonding nature and properties of compounds containing transition metals directly bound to a group 13 element is important partly because of their potential use as catalysts and in material science applications. In our search for synthetic methods leading to the formation of such compounds, we have reported the reaction between IrHCl₂(PCy₃)₂ and GaMe₃, and GaMe₂Cl with no evidence of heterobimetallic products. Under similar conditions, the analogous reaction of Vaska's gallyl compound (GaMe₃) to then proceed to a transition state \([\text{GaMe}_3\text{Ir(CO)(PPh}_3\text{)}_2]\) to finally release the alkylated product \(\text{trans-}[\text{GaMe}_2\text{Cl}]\) together with GaMe₂Cl.

The mechanism of methylation of Vaska's complex \(\text{trans-}[\text{ClIr(CO)(PPh}_3\text{)}_2]\) by trimethylgallium was studied and the identification of the spectroscopically detected intermediates was achieved with the aid of computational methods. The reaction pathway, computed by means of density functional theory (M05-2X-D3/def2-SVP), involves the initial formation of a chloride-bridged adduct \(\text{trans-}[\text{(Cl-GaMe}_3\text{)}\text{Ir(CO)(PPh}_3\text{)}_2]\) to then proceed to a transition state \(\text{[(Me}2\text{Cl-C-ClMeGaMe}_2\text{Ir(CO)(PPh}_3\text{)}_2]\). This transition state subsequently evolves to the methylated adduct \(\text{[MeIr(CO)(PPh}_3\text{)}_2\text{] (GaMe}_2\text{Cl})\) to finally release the alkylated product \(\text{trans-}[\text{MeIr(CO)(PPh}_3\text{)}_2\text{]}\) together with GaMe₂Cl.

Back to basics: identification of reaction intermediates in the mechanism of a classic ligand substitution reaction on Vaska's complex†

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The reported experimental investigations on ligand substitution reactions on complex 1 have mainly concentrated on the kinetics and associative and exchange mechanisms have been proposed for ligand substitutions. From a strict point of view, the elucidation of reaction mechanisms needs not only a good knowledge of the kinetics, which often gives valuable information, but also of the properties of the intermediates involved. In other words, how the reaction rate varies with the concentration of reactants, temperature, and other conditions gives vital clues but rarely establishes with certainty the mechanism of the process. In other words, how the reaction rate varies with the concentration of reactants, temperature, and other conditions gives vital clues but rarely establishes with certainty the mechanism of the process. Herein we use density functional theory (DFT) to assign the identity of the intermediates for the ligand substitution reaction of chloride by methyl in 1 which indeed proceeds via an associative mechanism involving a chloride-bridged transition state.
Materials and methods

General experimental procedures

All experiments were performed under argon atmosphere inside an MBraun glove box. Deuterated benzene was dried over molecular sieves and degassed via three freeze–pump–thaw cycles. Compound 1\(^{24}\) and authentic compound 2 were synthesized for comparison according to the reported procedures.\(^{25,26}\) NMR spectra were recorded on a Varian Inova 400 MHz in NMR tubes fitted with J. Young’s valves and using 0.5 mL solvent volume. IR spectra were collected on a Nicolet FTIR 6700 series spectrophotometer.

Computational details

All geometries were fully optimized using the M05-2X\(^{27}\) functional in conjunction with a def2-SVP basis set employing Gaussian 09 program package.\(^{28}\) In order to incorporate dispersion, Grimme’s D3 dispersion scheme\(^{29}\) was included during the optimization. The nature of stationary points located on the potential energy surface (PES) was characterized by harmonic vibrational frequency analysis. Thermal corrections were computed within the harmonic approximation. An intrinsic reaction coordinate (IRC) calculation was done in order to verify that the transition state was truly connected to the correct minima. Solvation energy corrections for benzene were evaluated by a self-consistent reaction field (SCRF) using the Solvation Model based on Density (SMD)\(^{30}\) model under the same level. All calculations were done using an ultrafine integration grid.

In order to evaluate the thermal effects, we used the procedure described by Irikura as is implemented in thermo code,\(^{31}\) where the estimation of the standard molar entropy and enthalpy change is computed from the molecular partition function. All the quantities needed are taken from the harmonic vibrational frequency computations.

The nature of the interactions was analyzed by energy decomposition analysis (EDA)\(^{32}\) at the PBE-D3/DZP/M05-2X-D3/def2-SVP level using the ADF (2013.01) package\(^{33}\) (TZVP basis set is used for Ir). Frozen core approximation was not used here. Scalar relativistic effects were considered using the zeroth-order regular approximation (ZORA).\(^{34}\) In EDA, bond formation between the interacting fragments can be represented by the following three successive steps. Firstly, the fragments, which are calculated with the frozen geometry of the entire molecule, are superimposed without electronic relaxation. It yields the quasiclassical electrostatic attraction (\(\Delta E_{\text{elstat}}\)). Secondely, the product wave function becomes antisymmetrized and renormalized providing the Pauli repulsion term (\(\Delta E_{\text{Pauli}}\)). In the next step the molecular orbitals undergo relaxation to go into their final form yielding the stabilizing orbital interaction (\(\Delta E_{\text{orb}}\)). Since we used dispersion corrected PBE-D3 functional, the dispersion correction term (\(\Delta E_{\text{disp}}\)) will be added to the interaction energy (\(\Delta E_{\text{int}}\)) values to describe the total bond energy as

\[
\Delta E_{\text{int}} = \Delta E_{\text{Pauli}} + \Delta E_{\text{elstat}} + \Delta E_{\text{orb}} + \Delta E_{\text{disp}}
\]

The \(\Delta E_{\text{int}}\) is related to the bond dissociation energy, \(D_{e}\), by adding \(\Delta E_{\text{prep}}\), which is the necessary energy needed to promote the fragments from their equilibrium geometry to the geometry in the complexes (eqn (4)). The advantage of using \(\Delta E_{\text{int}}\) instead of \(D_{e}\) is that the instantaneous electronic interaction of the fragments becomes analyzed, which yields a direct estimate of the energy components.

\[
-D_{e} = \Delta E_{\text{prep}} + \Delta E_{\text{int}} \quad (4)
\]

Results and discussion

Spectroscopic evidence

The reactions between 1 and GaMe\(_3\) in different stoichiometric ratios were monitored by NMR and IR spectroscopies. For spectroscopic studies, the addition of the organogallium was made upon frozen mixtures of 1 in 0.5 mL benzene-\(d_6\) and the spectra recorded as soon as possible after thawing. At the molar ratio 0.5 : 1 (corresponding to GaMe\(_3\) : 1), the NMR data were in agreement with partial formation of the methylated product 2.\(^{21}\) In addition, the \(^1\)H NMR spectrum recorded at 280 K shows two singlet signals of small intensity at \(\delta = 0.22\) and 0.08 ppm. The former signal is assigned to the generated GaMe\(_3\)Cl whilst the latter to intermediate species I\(_1\) (Table 2). Due to the low concentration of I\(_1\) and the fact that this species does not persist in solution for longer than 5 min at 280 K, it was not possible to record its \(^{31}\)P NMR spectrum. A sample prepared under identical conditions (benzene-\(d_6\)) was studied by IR spectroscopy in solution showing an additional carbonyl stretching band at 1965 cm\(^{-1}\), which is detectable only at short (<5 min) reaction times and we thus attributed it also to species I\(_1\) (Table 2). The NMR tube containing the reaction mixture GaMe\(_3\) : 1 (0.5 : 1 molar ratio) was allowed to warm up to 298 K during 5 min after mixing and the \(^1\)H NMR spectrum was recorded again. It showed evidence for the exchange of free and coordinated GaMe\(_3\) in the alkyl region as the signal due to the latter (at \(\delta = 0.15\) ppm at 280 K) and the new small signal at \(\delta = 0.08\) ppm merge into a broad signal in the region of \(-0.04 < \delta < 0.75\) ppm. The reaction of GaMe\(_3\) and 1 in a 1 : 1 molar ratio was performed under the same conditions and the \(^1\)H NMR spectrum collected immediately afterwards showed near-to-complete conversion of 1 to 2, as well as a new signal of small intensity at \(\delta = 0.38\) (t), which persists for few minutes later fully evolving to 2. Once again, rapid transformation to products and low concentration of the possible intermediate species prevented the acquisition of the \(^{31}\)P NMR spectrum. An independently but identically prepared sample for IR spectroscopy showed the presence of a new band at 1970 cm\(^{-1}\) (Table 2).

After comparison with reported frequencies of other adducts of 1,\(^{20,21,11,12}\) we propose that these experimental findings can be understood if a first intermediate species, I\(_1\), is formed at lower

\(\dagger\) The formation of 2 was previously known to occur from excess MeLi addition to 1 followed by addition of magnesium silicate, however no further details were provided for that transformation Ref. 3. The main spectroscopic features of an authentic sample of 2 prepared as reported in ref. 3 are a triplet at \(\delta 0.48\) ppm (\(J_{\text{HH}} = 9\) Hz) in the \(^1\)H NMR, a singlet at 31.9 ppm in the \(^{31}\)P\(^{(1)}\)H NMR in C\(_6\)D\(_6\) and a carbonyl stretching band at 1935 cm\(^{-1}\).
concentrations of GaMe₃, which incorporates at least one intact molecule of GaMe₃ such as to be able to exchange it with free GaMe₃ present in the reaction medium. The carbonyl band at 1965 cm⁻¹ in the IR spectrum and the ¹H signal at δ = 0.08 ppm would be then due to I₁. Addition of further organogallium (ratio 1 : 1 or higher) results in the formation of 2, together with species, I₂, observed by IR at 1970 cm⁻¹ and in the ¹H NMR at δ = 0.38 ppm as a triplet. The formation of GaMe₂Cl was spectroscopically ascertained by ¹H and ¹³C NMR spectra.

Investigation of the reaction mechanism by DFT

Since it was not possible by experimental means to ascertain the identity of the intermediate species, we turned to computational methods to gain insights into their structure. The electrostatic potential of 1, plotted in Fig. 1, clearly shows the presence of negative regions around Cl, O and Ir, indicating that these sites are susceptible for an electrophilic attack. According to the NBO charges, Ir exhibits a charge of −0.70|e|, followed by O (−0.49|e|) and Cl (−0.43|e|). So, compound 1 (Fig. 1) has three probable sites to coordinate a with GaMe₃ unit. In Fig. 2, these options are depicted: in (A) the galliumium coordinates through the chlorine atom, in (B) via the oxygen atom of the carbonyl group, and in (C) through the iridium itself.

Our M05-2X-D3/def2-SVP computations indicate that A is the lowest-lying energy structure. B and C are higher in energy than A by 4.1 and 0.5 kcal mol⁻¹, respectively. At 280 K (the temperature employed to record the ¹H NMR spectrum), the thermal effects change moderately the above energy differences (ΔG_A−B₂₈₀K = 3.0 kcal mol⁻¹ and ΔG_A−C₂₈₀K = 2.8 kcal mol⁻¹). Similar results are obtained using other functionals when dispersion is included (Table S1†). On the other hand, assuming a scale factor of 0.914 (details in the ESI†), the calculated frequencies for the CO stretching vibrations are 1964 (A), 1864 (B), and 1970 cm⁻¹ (C). The experimental value for 1I₁ is 1965 cm⁻¹ (Table 2). In principle, the chlorine-bridged adduct trans-[[Me₃Ga–Cl]Ir(CO)(PPh₃)₂] is the best candidate for 1I₁.

The temperature significantly affects the bond dissociation energies (BDE) of the three complexes. The computed dissociation energies at 0 K, including the ZPE corrections, are 20.4 (A), 16.3 (B), and 19.9 (C) kcal mol⁻¹. At 280 K, the calculated BDE values are 6.0 (A), 3.0 (B), and 3.2 (C) kcal mol⁻¹. Fig. 3 shows the variation of the free energy differences with temperature. It indicates that while A is stable until 470 K, B and C are stable complexes until only 320 K. Populations of each species estimated via Boltzmann distributions suggest that at 280 K and 1 atm, structures B and C contribute with less than 1%.

To gain a quantitative insight into the nature of the bonding interactions in these complexes, we also perform an EDA, using GaMe₃ and 1 as interacting fragments for 1I₁, B, and C, and GaMe₃Cl and 2 for I₂. Clearly, the principal factor to stabilize all four species is of electrostatic nature (around 46.5 ± 4.0%, Table 1). The orbital term contributes with around 37.0 ± 3.0%. The most important differences among the complexes come from the dispersion term. The dispersion term in complex B is more significant than in complexes A and C. It is well known that interactions with a high noncovalent nature are more affected by the temperature. Though in all of our complexes the dispersion is important, species B will be more affected by temperature.

Note the large difference between the BDE and the ΔF_int values. It denotes strong structural changes of the reactants
during the formation of the complexes. Our computations show that the structural changes in the GaMe₂, fragment are the most energetically expensive. In other words, the deformation of the Lewis acid not only influences dissociation energies, but also changes the electron distribution of the acid as it triggers changes in the gallium hybridization.

**I1** shows an almost undistorted Ir center in a square planar geometry in which the Ir–Cl bond length has been very slightly elongated (2.470 Å) with respect to the corresponding distance in the calculated unrestrained species 1 at 2.419 Å (Fig. 1). The Cl···Ga distance at 2.600 Å is still longer than that reported for covalent Ga–Cl bonds (ca. 2.23 Å)\(^\text{23,24}\) but it is close enough to establish a bonding interaction. This is presumably due to the fact that between the solid state structure and that in the gas phase for a dative bond the variations could be higher than 0.1 Å.\(^\text{25,26}\) In the extreme cases, these differences could be even higher than 0.8 Å. The Ir–Cl–Ga bond angle of 123.2° for the adduct **I1** is considerably more acute than that in the methylation of Vaska’s complex. Representative bond lengths (in Å) are shown, in parenthesis the corresponding lengths found upon inclusion of solvent. Atom colour code: C (black), Cl (light blue), Ga (olive green), H (grey), Ir (purple), O (red), P (orange).

![Fig. 4](image_url)

**Fig. 4** M05-2X-D3/def2-SVP geometry of the transition state involved in the methylation of Vaska’s complex. Representative bond lengths (in Å) are shown, in parenthesis the corresponding lengths found upon inclusion of solvent. Atom colour code: C (black), Cl (light blue), Ga (olive green), H (grey), Ir (purple), O (red), P (orange).

<table>
<thead>
<tr>
<th>Complex</th>
<th>(\Delta E_{\text{clat}})</th>
<th>(\Delta E_{\text{Pauli}})</th>
<th>(\Delta E_{\text{disp}})</th>
<th>(\Delta E_{\text{int}})</th>
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<tr>
<td><strong>I1</strong></td>
<td>–28.5 (45.3)</td>
<td>37.6</td>
<td>–23.5 (37.4)</td>
<td>–10.9 (17.3)</td>
</tr>
<tr>
<td><strong>Complex B</strong></td>
<td>–18.2 (43.1)</td>
<td>23.9</td>
<td>–14.3 (33.9)</td>
<td>9.7 (23.0)</td>
</tr>
<tr>
<td><strong>Complex C</strong></td>
<td>–40.1 (46.7)</td>
<td>51.9</td>
<td>–32.1 (37.4)</td>
<td>–13.6 (15.9)</td>
</tr>
<tr>
<td><strong>I2</strong></td>
<td>–73.5 (50.8)</td>
<td>93.3</td>
<td>–58.2 (40.2)</td>
<td>–13.0 (9.0)</td>
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</tbody>
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**Table 1** EDA results of the coordination of GaMe₂ to I1 studied at the PBE-D3/DZP//M05-2X-D3/dft2-SVP level. Energies are in kcal mol\(^{-1}\).

Further geometrical analysis of **TS** reveals a short calculated Ir–H distance of 2.097 Å between one of the hydrogens of the methyl group and the metal centre. This distance is short enough to consider it of agostic type. The agostic H···Ir bond angle (46.3°) is considerably smaller than those of the corresponding angles for the non-agostic hydrogens on the same carbon atom (120.5° and 125.8°). Other geometrical parameters such as the Ir–H···C angle of 111.0° and the H···C-Ga angle at 135.8° are well within reported values of systems for which an agostic interaction is established.\(^\text{39}\)

Evolution of the reaction leads to the formation of **I2** (see Fig. S1†) in which the methylated product has formed but which bears a GaClMe₂ fragment coordinated to Ir in the apical position of a square pyramid. At 0 K, **I2** is more stable than **I1** by 3.3 kcal mol\(^{-1}\). However, at 280 K the free energy difference is negligible, but the transformation from **I1** to **I2** is still exergonic (\(\Delta G_{\text{I1-I2}}^{298\text{K}} = -0.7\) kcal mol\(^{-1}\)). No agostic C···H···Ir interactions are established here and the structural parameters are rather close to the isolated final product. The computed Ir···Ga distance in **I2** of 2.607 Å is still larger than the other three X-ray structures reported to date, bearing direct Ir···Ga bonds, namely [M(15-COD)(IMes)(GaMe₂Cl)(Cl-GaMe₂)],\(^\text{2}^\) which has an Ir···Ga bond of 2.381(1)-2.389(2) Å, [Ir(GaMe₂Cl)CMe=CH₂][N(SiMe₂CH₂PPPPh₂)] of 2.411(2) Å,\(^\text{40}\) and [Ir(1,5-COD)(IMes)[Ga[N(4H)]Cl]₄] \([\text{IMes} = \text{C}[\text{N(C₆H₄Me₃}-\text{2,4,6}](\text{CH₃})₂] \) of 2.448(7) Å.\(^\text{41}\) Note that the Ir···Ga length of **I2** is however shorter than that of complex C (2.846 Å), indicating a weak heterobimetallic interaction. Our computations show that the dissociation of **I2** into **II** (Fig. S2†) and GaMe₂Cl (12.9 kcal mol\(^{-1}\)) is two times higher than that of **I1** (6.0 kcal mol\(^{-1}\)); so **I2** is considerably more stable than **I1**. The computed frequency for the CO stretching vibration of **I2** is 1984 cm\(^{-1}\), which is in good agreement with the experimental value of 1970 cm\(^{-1}\) (Table 2).

The computed mechanism is in agreement with other reports on ligand replacement or addition in low-valent transition metal complexes,\(^\text{42}\) aside from the proposed agostic interaction precluding the formation of the Ir–Me bond. However, this type of interaction has been well identified by others when counting for the stabilization of reaction intermediates or transient species\(^\text{43}\) in a variety of reaction mechanisms. Macgregor’s work on the stability of trans-(CO)Ir(XMe)(PPPh₃₂) (X = O, CH₃) is particularly relevant as agostic interactions are proposed to account for the transformation between subsequent transition states.\(^\text{44}\)

To this point the overall transformation from **I1** into **I2** is still somewhat unclear because the \(\Delta G_{\text{I1-I2}}^{298\text{K}}\) is 6.2 kcal mol\(^{-1}\).
Additionally, the structures of TS, I1 and I2, all possess large dipole moments (in the range of 5.2 to 7.7 D), and thus the geometries in the gas phase may be quite different from those in solution. However, since the reaction is performed using benzene as solvent we do not expect significant variations in the geometries and the energy profiles as a consequence of solvent effects (see ESI†). This latter point is evident from Fig. 1 and 2 in which the geometries do not show substantial changes when solvent is included in the computations. The largest variation is found in TS where the Ir–Cl bond distance changes by 0.04 Å. Moreover, Fig. 5 shows the free energy profiles for the reaction with and without solvent effects. It can be seen that inclusion of the solvent does not alter significantly the energy differences and most importantly the overall reaction is still endergonic. So, how to explain the formation of free compound 2?

Certainly, one possibility of a coupled reaction is the formation of gallium dimers. At the M05-2X-D3/def2-SVP level, the dimerization energy of GaMe3 at 280 K is endergonic (4.3 kcal mol⁻¹) in the gas phase. A similar value is obtained in solution (4.0 kcal mol⁻¹). In contrast, the dimerization energy of GaMe2Cl is −24.0 kcal mol⁻¹ in the gas phase at 280 K and −18.0 kcal mol⁻¹ in solution at the same temperature. These values suggest that dimerization of GaMe2Cl is the driving force for the conversion of 1 to 2. These results are in agreement with the reported structures of organogallium compounds both in solid state and in the gas phase. Indeed, in contrast to the monomeric nature of GaMe3,45 gallium chlorides Me3GaClx−y (x = 0, 1, 2) exhibit dimeric structures with a central Ga2Cl2 core.46 Our findings suggest that the aggregation of GaCl3Me as a dimer is crucial for the alkylation reaction from GaMe3.

Conclusions

The use of GaMe3 instead of more powerful alkylating ligands such as Li or Mg derivatives allows the spectroscopic detection of reaction intermediates under controlled conditions. The computational work has contributed to the assignment of the observed intermediate species I1 as trans-[Cl-GaMe3]Ir(CO)(PPh3)2 and I2 as [MeIr(CO)(PPh3)2]GaMe3. The transition state TS, comprising the two metal moieties in a metathesis-like reaction, is stabilized by an agostic C–H bond interaction. Although it is a simple ligand replacement reaction on a 16-electron Ir(i) complex, the present mechanism of alkylation of Vaska’s compound (1) by GaMe3 is in agreement with previous work highlighting the importance of the halide substituent on the reaction and the agostic interactions for the stabilisation of the transition state. In the present case, a GaMe3 adduct formed upon coordination to the chloride precedes the transition state. Finally, species TS and I2 both show weak interacting heterobimetallic Ir–Ga bonds.

Acknowledgements

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